

# Very Weak Solutions for Poisson-Nernst-Planck System

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Transport of Ionic Particles in Biological Environments

## **Outline**

**Introduction and Motivation**

**Very Weak Formulation**

**Elements of Proof**

**Applications to Poisson-Boltzmann and Other Systems**

**Conclusion**

## Poisson-Nernst-Planck System

$$\Omega \subset \mathbb{R}^n, n = 1, 2, 3$$

$$u_0, v_0 \in L^2_+(\Omega)$$

$$u, v \in L^\infty(0, T; L^2_+(\Omega)) \cap L^2(0, T; W^{1,2}(\Omega))$$

$$u_t = \nabla \cdot (\nabla u - u \nabla \phi)$$

$$v_t = \nabla \cdot (\nabla v + v \nabla \phi)$$

$$\Delta \phi = u - v, \quad x \in \Omega$$

$$(\nabla u - u \nabla \phi) \cdot \nu = (\nabla v + v \nabla \phi) \cdot \nu = 0$$

$$\phi = 0, \quad x \in \partial\Omega$$

## Navier-Stokes-Poisson-Nernst-Planck System

$$\Omega \subset \mathbb{R}^n, n = 1, 2, 3$$

$$\mathbf{u}_0 \in L^2(\Omega)^n, \nabla \cdot \mathbf{u}_0 = 0, \quad u_0, v_0 \in L_+^2(\Omega)$$

$$\mathbf{u} \in L^\infty(0, T; L^2(\Omega)^n) \cap L^2(0, T; W^{1,2}(\Omega)^n)$$

$$u, v \in L^\infty(0, T; L_+^2(\Omega)) \cap L^2(0, T; W^{1,2}(\Omega))$$

$$\mathbf{u}_t + \mathbf{u} \cdot \nabla \mathbf{u} + \nabla p = \Delta \mathbf{u} + \Delta \phi \nabla \phi, \quad \nabla \cdot \mathbf{u} = 0$$

$$u_t + \mathbf{u} \cdot \nabla u = \nabla \cdot (\nabla u - u \nabla \phi)$$

$$v_t + \mathbf{u} \cdot \nabla v = \nabla \cdot (\nabla v + v \nabla \phi)$$

$$\Delta \phi = u - v, \quad x \in \Omega$$

$$(\nabla u - u \nabla \phi) \cdot \nu = (\nabla v + v \nabla \phi) \cdot \nu = 0$$

$$\phi = 0, \mathbf{u} = 0, \quad x \in \partial\Omega$$

## Lyapunov Functional

$$\begin{aligned} W &= \int_{\Omega} \frac{|\mathbf{u}|^2}{2} + u \log u + v \log v \, dx - \int_{\Omega} \frac{|\nabla \phi|^2}{2} + \phi(u - v) \, dx \\ &= \int_{\Omega} \frac{|\mathbf{u}|^2}{2} + u \log u + v \log v \, dx + \int_{\Omega} \frac{|\nabla \phi|^2}{2} \, dx \end{aligned}$$

$$\frac{d}{dt} W = - \int_{\Omega} |\nabla \mathbf{u}|^2 + u |\nabla(\log u - \phi)|^2 + v |\nabla(\log v + \phi)|^2 \, dx$$

$$\mathbf{u} \in L^\infty(0, T; L^2(\Omega)^n) \cap L^2(0, T; W^{1,2}(\Omega)^n)$$

$$u, v \in L^\infty(0, T; L \log L(\Omega))$$

$$\phi \in L^\infty(0, T; W_0^{1,2}(\Omega))$$

## Energetic Variational Derivation

$$\begin{aligned} W &= \int_{\Omega} \frac{|\mathbf{u}|^2}{2} + u \log u + v \log v \, dx - \int_{\Omega} \frac{|\nabla \phi|^2}{2} + \phi(u - v) \, dx \\ &= \int_{\Omega} \frac{|\mathbf{u}|^2}{2} + u \log u + v \log v \, dx + \int_{\Omega} \frac{|\nabla \phi|^2}{2} \, dx \end{aligned}$$

$$-\frac{\delta W}{\delta \phi} = 0 \quad \Rightarrow \Delta \phi = u - v$$

$$-\frac{\delta^* W}{\delta^* u} = x_t^1 \quad \Rightarrow J = \nabla u - u \nabla \phi$$

$$-\frac{\delta^* W}{\delta^* v} = y_t \quad \Rightarrow K = \nabla v + v \nabla \phi$$

$$-\frac{\delta W}{\delta x} = F \quad \Rightarrow F = (u - v) \nabla \phi = \nabla \cdot (\nabla \phi \otimes \nabla \phi) - \nabla \frac{1}{2} |\nabla \phi|^2$$

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<sup>1</sup> $\delta u + \nabla \cdot (u \omega) = 0$

## Mathematical Analysis in 2 and 3 Dimensions

### Theorem (Biler, et al 1994)

For  $\Omega \in \mathbb{R}^2$  and  $u_0, v_0 \in L_+^2$  there exists a unique solution  $u, v \in L^\infty(0, T; L_+^2) \cap L^2(0, T; W^{1,2})$  of the PNP system and for  $\tau > 0$ ,  $u, v \in C^{2,\alpha}((\tau, T) \times \Omega)$ .

$$\|w\|_{L^3}^3 \leq \epsilon \|\nabla w\|_{L^2}^{3/2} + (1 + \|w\|_{L^1}^{3/2}) \exp(C\|w\|_{L \log L} \epsilon^{-1})$$

### Theorem (Biler, et al 1994)

For  $\Omega \in \mathbb{R}^3$  and  $u_0, v_0 \in L_+^2$  there exists a unique solution  $u, v \in L^\infty(0, T_0; L_+^2) \cap L^2(0, T_0; W^{1,2})$  of the PNP system and for  $\tau > 0$ ,  $u, v \in C^{2,\alpha}((\tau, T_0) \times \Omega)$ .

$$T_0 = T_0(\Omega, \|u_0, v_0\|_{L^2})$$

### Conjecture

The maximal interval of existence is  $[0, \infty)$ .

## Recent Mathematical Analysis in 2 and 3 Dimensions

- ▶ Debye system: Biler, Hebisch, Hilhorst, Nadzieja (1992-1994)
- ▶ Asymptotics and long time behavior: Biler, Dolbeault (2000)
- ▶ Hydrodynamic coupling: Jerome, Sacco (2009), RJR (2009), Schmuck (2009), Deng, Zhao, Cui (2010), Jerome (2011)
- ▶ Diffusion limit of Vlasov-Poisson-Fokker-Planck system: Wu, Lin, Liu (2014)
- ▶ Quasilinear approximation with capacitance: Bothe, Fischer, Pierre, Rolland (2014)

## Motivation

Lyapunov functional for a variety of boundary conditions places

ionic densities uniformly in  $L^\infty(0, T; L^1(\Omega))$

potential in  $L^\infty(0, T; W^{1,2}(\Omega))$

- \* Pose a very weak formulation with regularity informed by the a priori estimates – infer regularity from very weak formulation \*

## Very Weak Formulation : Derivation

Total charge system resembles Debye system for single anion.

$$u_t - \nabla \cdot (\nabla u - u \nabla \phi) = f$$

$$\Delta \phi = u$$

Relationship between convection and Maxwell stress:

$$-u \nabla \phi = -\Delta \phi \nabla \phi = \nabla \frac{1}{2} |\nabla \phi|^2 - \nabla \cdot (\nabla \phi \otimes \nabla \phi)$$

$$u_t - \Delta(u + \frac{1}{2} |\nabla \phi|^2) + \nabla \cdot \nabla \cdot (\nabla \phi \otimes \nabla \phi) = f$$
$$\Delta \phi = u$$

## Very Weak Formulation

$$\boxed{\begin{aligned} u_t - \Delta(u + \tfrac{1}{2}|\nabla\phi|^2) + \nabla \cdot \nabla \cdot (\nabla\phi \otimes \nabla\phi) &= f \\ \Delta\phi &= u \end{aligned}}$$

$$u \in L^\infty(0, T; L^1(\Omega)), \quad \phi \in L^\infty(0, T; W^{1,2}(\Omega))$$

$$\begin{aligned} & \int_0^T \int_{\Omega} -uw_t - (u + \tfrac{1}{2}|\nabla\phi|^2)\Delta w + \nabla\phi \otimes \nabla\phi : \nabla^2 w \, dx \, dt \\ &= \int_0^T \int_{\Omega} fw \, dx \, dt \\ & \int_0^T \int_{\Omega} -\nabla\phi \cdot \nabla\eta \, dx \, dt = \int_0^T \int_{\Omega} u\eta \, dx \, dt \end{aligned}$$

## Very Weak Solutions in Other Contexts

- ▶ Weyl's lemma

$$u \in L^1(\mathbb{R}^n), \quad \int_{\mathbb{R}^n} u \Delta w \, dx = 0, \quad \forall w \in C_0^\infty(\mathbb{R}^n)$$
$$\Rightarrow u \text{ analytic, } \Delta u = 0$$

- ▶ Semi-linear 2nd order equations in  $L^1$  : Brezis, Strauss (1973), Brezis, Cazenave, Martel (1996)
- ▶  $p$ -Laplacian and heat flow thereof: Lewis (1993), Kinnunen, Lewis (2002)
- ▶ Semi-linear elliptic equations in boundary weighted spaces: Diaz, Rakotoson (2009, 2010)
- ▶ Divergence form elliptic equations: Jin, Mazya, Schaftingen (2009), Zhang, Bao (2012)
- ▶ Coarea flow: Hardt, Zhou (1994), Hardt, Tonegawa (1996)

## Very Weak Formulation : Stationary Problem

$$\begin{aligned}\Delta(u + \frac{1}{2}|\nabla\phi|^2) - \nabla \cdot \nabla \cdot (\nabla\phi \otimes \nabla\phi) &= -f \\ |\Delta\phi| &\leq u\end{aligned}$$

$$u \in L^1_{\text{loc}}(\Omega), \quad \phi \in W^{1,2}_{\text{loc}}(\Omega)$$

$$\begin{aligned}\int_{\Omega} (u + \frac{1}{2}|\nabla\phi|^2) \Delta w - \nabla\phi \otimes \nabla\phi : \nabla^2 w \, dx &= - \int_{\Omega} fw \, dx \\ \int_{\Omega} -\nabla\phi \cdot \nabla\eta = \int_{\Omega} \xi\eta \, dx, \quad |\xi|(x) &\leq u(x)\end{aligned}$$

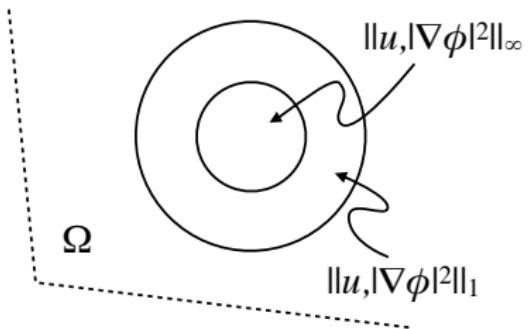
## Main Result : Interior Regularity

### Theorem (1)

If  $(u, \phi) \in L^1_{\text{loc},+}(\Omega) \times W^{1,2}_{\text{loc}}(\Omega)$  is a very weak solution of

$$\Delta(u + \frac{1}{2}|\nabla\phi|^2) - \nabla \cdot \nabla \cdot (\nabla\phi \otimes \nabla\phi) = -f, \quad |\Delta\phi| \leq u$$

with  $\Delta^{-1}f \in L^\infty(\Omega)$ , then  $u, \nabla\phi \in L^\infty_{\text{loc}}(\Omega)$ .



## Interior Regularity : Multiple Species

### Theorem (2)

If  $(u_1, \dots, u_q, \phi) \in (L_{\text{loc},+}^2(\Omega))^q \times W_{\text{loc}}^{1,2}(\Omega)$  is a **weak solution** of

$$\nabla \cdot (D_1(x)\nabla u_1 - z_1 u_1 \nabla \phi) = -f_1$$

⋮

$$\nabla \cdot (D_q(x)\nabla u_q - z_q u_q \nabla \phi) = -f_q$$

$$\Delta \phi = z_1 u_1 + \dots + z_q u_q + g$$

then  $u_1, u_2, \dots, u_q, \nabla \phi \in C^\infty(\Omega)$ .

## Elements of Proof : Monotonicity Formula

If  $f \in L^1_{\text{loc}}(\Omega)$ , then for almost every  $x \in \Omega$ ,

$$\lim_{\rho \rightarrow 0} \frac{1}{\rho^n} \int_{B_\rho(x)} f(z) dz = f(x)|B_1(0)|$$

and

$$\lim_{\rho \rightarrow 0} \frac{1}{\rho^p} \int_{B_\rho(x)} f(z) dz = 0, \quad p < n$$

## Elements of Proof : Monotonicity Formula

$$\nabla \cdot (\nabla u + \frac{1}{2} |\nabla \phi|^2) - \nabla \cdot \nabla \cdot (\nabla \phi \otimes \nabla \phi) = -f$$

↓

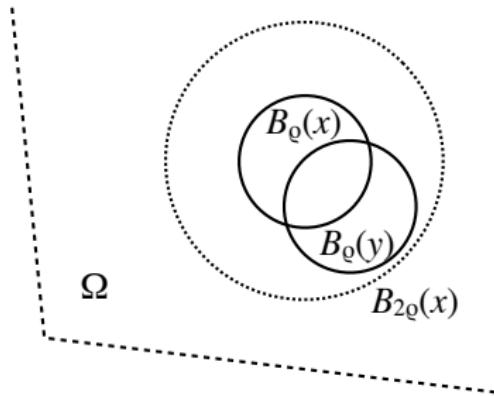
$$\nabla \cdot (\nabla u_\epsilon + \frac{1}{2} |\nabla \phi|_\epsilon^2) - \nabla \cdot \nabla \cdot (\nabla \phi \otimes \nabla \phi)_\epsilon = -f_\epsilon$$

↓

$$\begin{aligned} & \int_{B_r(x)} (2 - \beta)u \, dz - r \int_{\partial B_r(x)} (\nabla \phi \cdot \nu)^2 \, dS = r^{n-1+\beta} \frac{d}{dr} \frac{P_{x,r}}{r^{n-2+\beta}} \\ & + \int_{B_r(x)} \frac{\beta}{2} |\nabla \phi|^2 \, dz + \int_{B_r(x)} f \frac{1}{2} (r^2 - |z - x|^2) \, dz \end{aligned}$$

$$P_{x,r} = \int_{B_r(x)} u + \frac{1}{2} |\nabla \phi|^2 \, dz, \quad U_{x,\beta,\rho} = \int_0^\rho \frac{1}{r^{n-1+\beta}} \int_{B_r(x)} u \, dz \, dr$$

## Riesz-like Function Space



$$U_{x,\beta,\rho} = \int_0^\rho \frac{1}{r^{n-1+\beta}} \int_{B_r(x)} u \, dz \, dr, \quad \mathcal{U}_{x,\beta,\rho} = \operatorname{ess\,sup}_{y \in B_\rho(x)} U_{y,\beta,\rho}$$

$$L^\infty \subsetneq \{u : \mathcal{U}_{x,\beta,\rho} < \infty\} \subsetneq L^1, \quad \beta \in [0, 2)$$

$$M^p \subsetneq \{u : \mathcal{U}_{x,\beta,\rho} < \infty\}, \quad \beta \in [0, 2 - n/p)$$

## Riesz-like Function Space

$$(n - 2 + \beta)U_{x,\beta,\rho} + \frac{1}{\rho^{n-2+\beta}} \int_{B_\rho(x)} u \, dz = \int_{B_\rho(x)} \frac{u(z)}{|z - x|^{n-2+\beta}} \, dz$$

$$\begin{aligned} |\nabla^k \phi(x) - \nabla^k \phi(y)| &\leq \int_{B_{r/2}(x)} u(z) |K_k(z, x) - K_k(z, y)| \, dz \\ &+ \int_{B_{r/2}(y)} u(z) |K_k(z, x) - K_k(z, y)| \, dz \\ &+ \int_{B_{2\rho}(x_0) \setminus (B_{r/2}(x) \cup B_{r/2}(y))} u(x) |K_k(z, x) - K_k(z, y)| \, dz \\ &+ \int_{B_{2\rho}(x_0)} |\phi|(x) |L_k(z, y) - L_k(z, y)| \, dz \end{aligned}$$

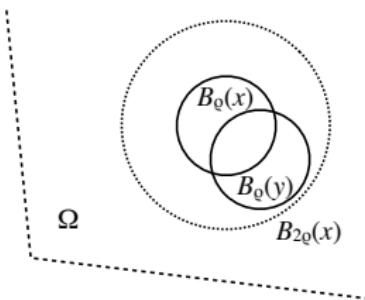
# Riesz-like Function Space

## Proposition

Let  $\beta \in (0, 1)$  and  $\phi \in W_{\text{loc}}^{1,2}(\Omega)$  satisfy  $|\Delta \phi| \leq u$

1. If  $\mathcal{U}_{x,0,\rho} < \infty$ , then  $\phi \in L^\infty(B_\rho(x))$
2. If  $\mathcal{U}_{x,\beta,\rho} < \infty$ , then  $\phi \in C^{0,\beta}(B_\rho(x))$
3. If  $\mathcal{U}_{x,1+\beta,\rho} < \infty$ , then  $\phi \in C^{1,\beta}(B_\rho(x))$

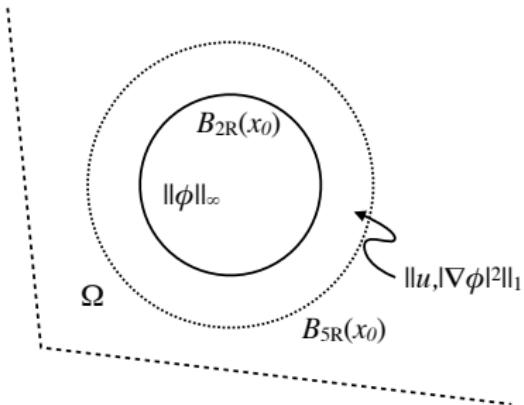
$$U_{x,\beta,\rho} = \int_0^\rho \frac{1}{r^{n-1+\beta}} \int_{B_r(x)} u \, dz \, dr, \quad \mathcal{U}_{x,\beta,\rho} = \operatorname{ess\,sup}_{y \in B_\rho(x)} U_{y,\beta,\rho}$$



## Elements of Proof : Zeroth Order Bounds

$$\mathcal{U}_{x_0,0,\rho} < \infty \text{ and } \phi \in L_{\text{loc}}^{\infty}(\Omega)$$

$$\begin{aligned} \int_{B_r(x)} 2u \, dz &= r^{n-1} \frac{d}{dr} \frac{P_{x,r}}{r^{n-2}} - r \int_{\partial B_r(x)} (\nabla \phi \cdot \nu)^2 \, dS \\ &+ \int_{B_r(x)} f \frac{1}{2} (r^2 - |z - x|^2) \, dz \end{aligned}$$



## Trace Inequality : Bounded to Hölder Continuity

$\phi \in L^\infty_{\text{loc}}(\Omega)$  and  $|\Delta\phi| \leq u$

$$\int_{B_r(x)} |\nabla\phi|^2 dz \leq C_0 \left( r \int_{\partial B_r(x)} (\nabla\phi \cdot \nu)^2 dS + \int_{B_r(x)} u dz \right)$$

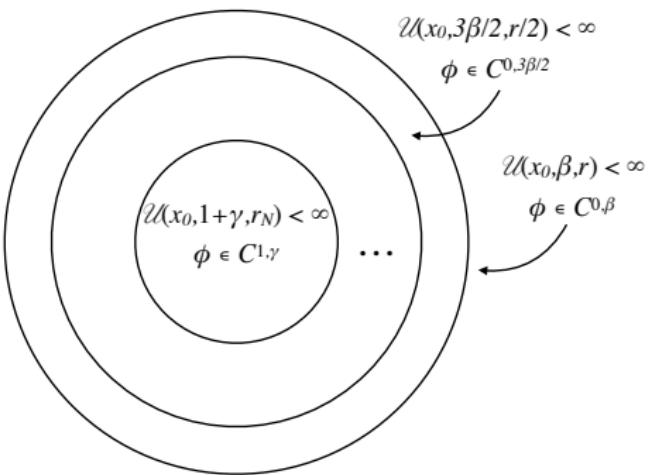
⇓

$$\begin{aligned} (2 - \beta - C\beta) \int_{B_r(x)} u dz &= r^{n-1+\beta} \frac{d}{dr} \frac{P_{x,r}}{r^{n-2+\beta}} \\ &+ \int_{B_r(x)} f^{\frac{1}{2}}(r^2 - |z-x|^2) dz \end{aligned}$$

## Sequential Improvement

$$\mathcal{U}_{x_0, \beta, \rho} < \infty \text{ and } \phi \in C_{\text{loc}}^{0, \beta}(\Omega)$$

$$\int_{B_r(x)} |\nabla \phi|^2 dz \leq \frac{2r}{\beta_{i+1}} \int_{\partial B_r(x)} (\nabla \phi \cdot \nu)^2 dS + C \left( \boxed{r^{n-2+2\beta}} + r^\beta \int_{B_r(x)} u dz \right)$$



## Bounding the Excess

Monotonicity formula when  $\beta = 2$ :

$$0 \leq \frac{d}{dr} \frac{P_{x,r}}{r^n} + \frac{E_{x,r}}{r^{n+1}} + \frac{1}{r^{n+1}} \int_{B_r(x)} f^{\frac{1}{2}}(r^2 - |z-x|^2) dz.$$

Mean value excess:

$$E_{x,r} = \int_{B_r(x)} |\nabla \phi|^2 dz - r \int_{\partial B_r(x)} (\nabla \phi \cdot \nu)^2 dS \leq Cr^{n+\gamma}$$

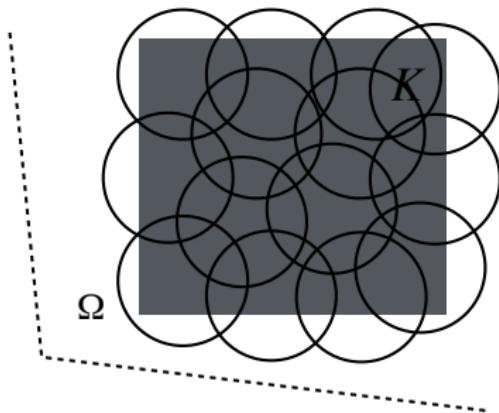
$$\begin{aligned} & (u(x) + \frac{1}{2} |\nabla \phi|^2(x)) |B_1(0)| \\ & \leq \frac{P_{x,\rho}}{\rho^d} + \int_0^\rho Cr^{-1+\gamma} + \frac{1}{r^{n-1}} \int_{B_r(x)} |f| dz dr \end{aligned}$$

## Inner Estimate from Outer Bound

### Lemma

For  $x_0 \in \Omega$  and  $B_{5r_0}(x_0) \subset \Omega$ , the number of levels  $N$  satisfies

$$N \leq C_1 + 3 \ln \left( C_T + P_{x_0, 4r_0} + \|\phi\|_{L^1(B_{4r_0}(x_0))} + \mathcal{F}_{x_0, 0, 2r_0} \right)$$



## Summary of Theorem and Proof

### Theorem (1)

If  $(u, \phi) \in L^1_{\text{loc},+}(\Omega) \times W^{1,2}_{\text{loc}}(\Omega)$  is a very weak solution of

$$\Delta(u + \frac{1}{2}|\nabla\phi|^2) - \nabla \cdot \nabla \cdot (\nabla\phi \otimes \nabla\phi) = -f, \quad |\Delta\phi| \leq u$$

with  $\Delta^{-1}f \in L^\infty(\Omega)$ , then  $u, \nabla\phi \in L^\infty_{\text{loc}}(\Omega)$ .

- ▶ Prove initial uniform and Hölder estimate
- ▶ Show sequential improvement of Hölder estimate
- ▶ Continue until gradient Hölder estimate reached
- ▶ Estimate mean value excess

## Application I : Multiple Species

$$\nabla \cdot (D_1(x) \nabla u_1 - z_1 u_1 \nabla \phi) = 0$$

⋮

$$\nabla \cdot (D_q(x) \nabla u_q - z_q u_q \nabla \phi) = 0$$

$$\Delta \phi = z_1 u_1 + \cdots + z_q u_q$$

$$u = u_1 + u_2 + \cdots + u_q$$

$$\Delta(u + \frac{1}{2}|\nabla \phi|^2) - \nabla \cdot \nabla \cdot (\nabla \phi \otimes \nabla \phi) = \sum_{i=1}^q \nabla \cdot ((1 - D_i) \nabla u_i)$$

$$|\Delta \phi| = \left| \sum_{i=1}^q z_i u_i \right| \leq u$$

## Application I : Multiple Species

$$f := - \sum_{i=1}^q \nabla \cdot ((1 - D_i) \nabla u_i)$$

$$\begin{aligned} & \int_0^\rho \frac{1}{r^{n-1+\beta}} \int_{B_r(x)} f \frac{1}{2} (r^2 - |z-x|^2) dz dr \\ & \leq \left( \max_{i=1,\dots,q} \|D_i\|_{C^0(B_\rho(x))} + 1 \right) \frac{1}{\rho^{n-2+\beta}} \int_{B_\rho(x)} u dz dr \\ & + \max_{i=1,\dots,q} \|\nabla D_i\|_{C^0(B_\rho(x))} U_{x_0, \beta-1, \rho} \end{aligned}$$

### Proposition

If  $(u_1, \dots, u_q, \phi) \in (L_{\text{loc},+}^2(\Omega))^q \times W_{\text{loc}}^{1,2}(\Omega)$  is a weak solution,  
then  $u_1, u_2, \dots, u_q, \nabla \phi \in L_{\text{loc}}^\infty(\Omega)$ .

## Further Applications

### Poisson-Boltzmann with Natural Charge Boundary Conditions

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$$\Delta\phi = \sum_{i=1}^q z_i e^{\lambda_i \phi_i}, \quad A\phi + B\nabla\phi \cdot \nu = \phi_0$$

### Poisson-Boltzmann with Fixed Charge Boundary Conditions

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$$\nabla \cdot (\nabla u_i - z_i u_i \nabla \phi) = f_i, \quad u_i = g_i$$

$$\Delta\phi = \sum_{i=1}^q z_i u_i, \quad A\phi + B\nabla\phi \cdot \nu = \phi_0$$

## Application II : Keller-Segel

### Theorem (3)

If  $(u, \phi) \in L^1_{\text{loc},+}(\Omega) \times W^{1,2}_{\text{loc}}(\Omega)$  is a very weak solution of

$$\nabla \cdot (\nabla u + u \nabla \phi) = 0, \quad \Delta \phi = u$$

with  $\Omega \subset \mathbb{R}^2$ , then  $u, \nabla \phi \in C^\infty_{\text{loc}}(\Omega)$ .

$$\int_{B_r(x)} n u - \frac{n-2}{2} |\nabla \phi|^2 dz = r \int_{\partial B_r(x)} u - \frac{1}{2} |\nabla \phi|^2 + (\nabla \phi \cdot \nu)^2 dS.$$

## Application II : Keller-Segel

### Theorem (3)

If  $(u, \phi) \in L^1_{\text{loc},+}(\Omega) \times W^{1,2}_{\text{loc}}(\Omega)$  is a very weak solution of

$$\nabla \cdot (\nabla u + u \nabla \phi) = 0, \quad \Delta \phi = u$$

with  $\Omega \subset \mathbb{R}^2$ , then  $u, \nabla \phi \in C^\infty_{\text{loc}}(\Omega)$ .

### Corollary

Stationary, two dimensional chemotaxis is singular measure valued or null (existence not claimed).

## Application II : Keller-Segel

### Theorem (3)

If  $(u, \phi) \in L^1_{\text{loc},+}(\Omega) \times W^{1,2}_{\text{loc}}(\Omega)$  is a very weak solution of

$$\nabla \cdot (\nabla u + u \nabla \phi) = 0, \quad \Delta \phi = u$$

with  $\Omega \subset \mathbb{R}^2$ , then  $u, \nabla \phi \in C^\infty_{\text{loc}}(\Omega)$ .

Counter example:  $\Omega \subset \mathbb{R}^{n \geq 3}$ :

$$u = 2(n-2)|x|^{-2} \in L^1_{\text{loc},+}(\Omega), \quad \phi = 2 \log |x| \in W^{1,2}_{\text{loc}}(\Omega).$$

is a very weak solution of

$$\nabla \cdot (\nabla u + u \nabla \phi) = 0, \quad \Delta \phi = u$$

## Time Dependent Formulation

$$\begin{aligned} & \int_0^T \int_{\Omega} -uw_t - (u + \frac{1}{2}|\nabla\phi|^2)\Delta w + \nabla\phi \otimes \nabla\phi : \nabla^2 w \, dx \, dt \\ &= \int_0^T \int_{\Omega} fw \, dx \, dt \\ & \int_0^T \int_{\Omega} -\nabla\phi \cdot \nabla\eta \, dx \, dt = \int_0^T \int_{\Omega} \xi\eta \, dx \, dt, \quad |\xi| \leq u \end{aligned}$$

### Conjecture

*A very weak solution*

$$u \in L^\infty(0, T; L^1(\Omega)), \quad \phi \in L^\infty(0, T; W^{1,2}(\Omega))$$

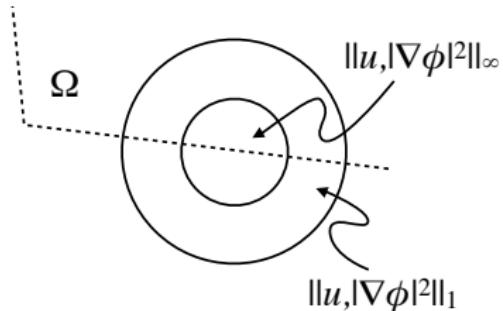
*enjoys*  $u, |\nabla\phi|^2 \in L_{\text{loc}}^\infty((0, T) \times \Omega)$ .

## Summary

- ▶ Viewed as Maxwell stress, total charge convection converted into 2nd order term
- ▶ Formulated weak system with solutions in the space  $L^1(\Omega) \times W^{1,2}(\Omega)$  appropriate to Lyapunov functional
- ▶ Proved interior  $L^\infty$  estimate based on  $L^1$  data
- ▶ Inferred full interior regularity for Poisson-Boltzmann with various boundary conditions
- ▶ Applied argument to two dimensional, stationary chemotaxis
- ▶ Further applications to regularization and compactness

## Conclusion

- ▶ Boundary regularity



- ▶ Variable dielectric  $\nabla \cdot (\epsilon(x) \nabla \phi) = \sum_{i=1}^q z_i u_i + g$

$$\begin{aligned}\underline{\epsilon} \int_{B_r(x)} |\nabla \phi|^2 dz &\leq \int_{B_r(x)} \epsilon(x) |\nabla \phi|^2 dz \\ &\leq \int_{\partial B_r(x)} \epsilon(x) |\phi - \bar{\phi}| |\nabla \phi \cdot \nu| dS \\ &\quad + \int_{B_r(x)} (u + |g|) |\phi - \bar{\phi}| dz\end{aligned}$$

## Conclusion

- ▶ Challenges :
  - derive mean value type monotonicity formula for time dependent problem – develop interior/boundary estimates
- ▶ Local monotonicity formulas in other contexts
  - ▶ Heat equation: Fulks (1971)
  - ▶ Semilinear diffusion equation: Giga, Kohn (1985)
  - ▶ Harmonic maps: Struwe (1988), Chen, Struwe (1989)
  - ▶ Mean curvature flow: Hamilton (1993), Ecker (2001)
- ▶ Thank you Chun and Maximilian!